

# Study Of Electron Transport Coefficients And Critical Field Strength In N<sub>2</sub>O And N<sub>2</sub>O-SF<sub>6</sub> Mixtures Using Boltzmann Equation Analysis

Mohammad Mustafa Othman<sup>1</sup>, Gullala YasenBakr<sup>2</sup>, Sherzad Aziz Taha<sup>3</sup>

<sup>1,2,3</sup>*Department of Physics, College of education, University of Salahaddin-Hawler, KRG, Iraq*

<sup>1</sup>*Department of Library Science, Tishik International university, KRG, Erbil*

*Email id: <sup>1</sup>muhamad.othman@su.edu.krd*

## **Abstract**

*equation over the density-normalized electric field strength,  $E/N$  from 0.1 to 1000 Td ( $1Td=10^{-17} V.cm^2$ ). The swarm parameters such as, drift velocity, electron mean energy, characteristic energy, ionization and attachment coefficients were analyzed using a set of cross sections for the gases. The density-reduced effective ionization coefficient  $(\alpha - \eta)/N$  has also been calculated. The present results agree well with previous experimental and theoretical results. Moreover, the  $E/N$  value at which ionization coefficient equals attachment coefficient, the critical field strength was calculated to be  $(E/N)_{cr}=199$  Td for pure N<sub>2</sub>O gas, which is lower than that of pure SF<sub>6</sub> gas (361 Td), a widely used as insulating gas. In addition, it's found that the critical field strength (voltage breakdown) increase with increasing SF<sub>6</sub> concentration in the gas mixtures.*

**Keywords:** Nitrous oxide, swarm parameters, cross section, critical field strength, breakdown voltage, EEDF.

## **1. INTRODUCTION**

Nitrous oxide (N<sub>2</sub>O) is an electronegative gas with breakdown voltage about half that of SF<sub>6</sub> (Biasiutti, 2013), has excellent insulation widely used in gas insulated transmission line and gas insulated switchgear. Nitrous oxide known as laughing gas with a low global warming potential after carbon dioxide and methane. N<sub>2</sub>O is atmospheric greenhouse gas with high global warming potential (GWP) is 310 times higher than that of CO<sub>2</sub> (as a reference, the GWP of CO<sub>2</sub>=1) and the life time is 120 years in the atmosphere on 100 year horizon (Gunnar Myhre et al., 2013). The concentration of nitrous oxide reached (333 ppb) in 2020 is much lower than that of CO<sub>2</sub> and CH<sub>4</sub>.

Nitrous oxide is a gas at room temperature, is 1.53 times denser than air, 34 times more soluble than nitrogen in blood, is colorless, tasteless, non-flammable, odorless, non-toxic and ozone depleting substance. N<sub>2</sub>O has stable chemical structures has a low boiling point of -88.5 C, and strong electrical insulation gas use in numerous applications, chemistry, medicine, media gas laser, plasma etching, high-voltage equipment, thin film by using plasma enhanced chemical vapour deposition and technology. Besides the GWP of N<sub>2</sub>O is lower than that of SF<sub>6</sub>, it has been suggested as a potential replacement to SF<sub>6</sub> gas which used in the electrical industry with a higher global warming potential 23900 times higher than that of CO<sub>2</sub>, and with long life time in atmosphere 3400 years (Kim, et al., 2004), and would allow

application at higher pressures than that SF<sub>6</sub>. In 1997 the Kyoto protocol has listed SF<sub>6</sub>, CO<sub>2</sub>, CH<sub>4</sub>, N<sub>2</sub>O and hydrofluorocarbon as strong greenhouse gases, and for environment required to reduce the use of this gases.

Electron swarm parameters is described by the following parameters: the electron drift velocity  $v_d$ , mean electron energy  $\langle \varepsilon \rangle$ , characteristic energy  $D/\mu_e$  (where  $D$  is diffusion coefficient and  $\mu_e$  electron mobility), reduced ionization coefficient  $\alpha/N$ , reduced attachment coefficient  $\eta/N$ . These swarm parameters have been measured in pure N<sub>2</sub>O for small values of the reduced electric field  $E/N < 490$  Td ( $1\text{Td}=10^{-17}$  V cm<sup>2</sup>), (Mechlinska-Drewko, 2003).

The drift velocity in pure N<sub>2</sub>O was measured by (Lisovskiy, et al., 2006) in strong electric field from RF breakdown curves in the range  $E/P$  varying from 87 to 840 Vcm<sup>-1</sup>Torr<sup>-1</sup>. The electron drift velocity and effective ionization coefficients in N<sub>2</sub>O, N<sub>2</sub>O-N<sub>2</sub> and N<sub>2</sub>O-SF<sub>6</sub> mixtures are measured by (Basurto, et al., 2007) over  $E/N$  varying from 0.5 to 400 Td. (Dupljanin, et al., 2010) measured the electron transport coefficients and electron collision cross sections in N<sub>2</sub>O and in N<sub>2</sub>O-N<sub>2</sub> mixtures. Recently the electron swarm coefficients and the limiting field strength of binary gas mixtures SF<sub>6</sub>-N<sub>2</sub>O measured by (Basurto, et al., 2013) with a pulsed Townsend technique over a wide range of  $E/N$  from 130 to 420 Td. Furthermore, one recent review (Andreas, et al., 2020) states that still relatively little is known on electron swarm parameters and ionization and attachment properties, and electrical insulation properties for this molecule. Finally, (Hainan Liu, 2020), studied the electron scattering with N<sub>2</sub>O molecule for control and reduction of atmospheric pollution.

In the present work we calculated the electron swarm parameters, effective ionization coefficients and critical field strength for N<sub>2</sub>O, SF<sub>6</sub> and different mixtures (5%, 20% and 80%) of SF<sub>6</sub>, by using two-term solution of the Boltzmann equation method, over a wide range of the reduced electric field strength  $E/N$  from 0.1 to 1000 Td, where  $E$  is the electric field and  $N$  is the gas density.

## 2. THEORETICAL FORMULATION

In this section we summarize the theoretical expression actually used in our calculation.

The electron energy distribution function  $f(\varepsilon)$  is obtained from the two-term solution of the Boltzmann equation (Colonna & D'Angola, 2016). The electron energy distribution function can be normalized.

$$\int_0^{\infty} f(\varepsilon) \sqrt{\varepsilon} d\varepsilon = 1 \quad (1)$$

The Boltzmann equation for electrons moves due to an external (dc) electric field  $\bar{E}$  is written as (Smith and Thomson, 1974; Phelps and Pitchford, 1982; Othman, et al., 2019),

$$\frac{\partial f(\mathbf{v})}{\partial t} + \mathbf{v} \cdot \nabla_r f(\mathbf{v}) - \frac{e}{m} \bar{E} \cdot \nabla_v f(\mathbf{v}) = \left( \frac{\partial f}{\partial t} \right)_{oll} \quad (2)$$

Where,  $e$  is electron charge,  $m$  is the electron mass, the right part of the equation denotes the rate of change in the electron distribution due to elastic and inelastic collisions. To solve the Boltzmann transport equation the electron distribution function is expanded in two terms of Legendre polynomials,  $f(\bar{\mathbf{v}}) = f_o(\mathbf{v}) + f_1(\mathbf{v}) \cos \theta$ , where  $f_o(\mathbf{v})$  is the isotropic and  $f_1(\mathbf{v})$  is anisotropic part of the distribution function where  $f_1(\mathbf{v}) \ll f_o(\mathbf{v})$ . The expansions in

spherical harmonic using two term solution have been given elsewhere (Govinda-Raju, G. 2006).

The electron swarm parameters are expressed in term of the electron energy distribution function  $f(\varepsilon)$  and collision cross sections  $Q_m^T(\varepsilon)$  as follows

The electron drift velocity is,

$$v_d = -\frac{\bar{E}}{3} \sqrt{\frac{2e}{m}} \int_0^\infty \frac{\varepsilon}{NQ_m^T(\varepsilon)} \frac{\partial f(\varepsilon)}{\partial \varepsilon} d\varepsilon \quad (3)$$

The electron mean energy is,

$$\langle \varepsilon \rangle = \int_0^\infty \varepsilon^{3/2} f(\varepsilon) d\varepsilon \quad (4)$$

The reduced transverse diffusion coefficient is,

$$D_T N = \frac{1}{3} \sqrt{\frac{2e}{m}} \int_0^\infty \frac{\varepsilon}{Q_m^T(\varepsilon)} f(\varepsilon) d\varepsilon \quad (5)$$

The reduced electron mobility is,

$$\mu_e N = -\frac{1}{3} \sqrt{\frac{2e}{m}} \int_0^\infty \frac{\varepsilon}{Q_m^T(\varepsilon)} \frac{\partial f(\varepsilon)}{\partial \varepsilon} d\varepsilon \quad (6)$$

The characteristic energy is,

$$\varepsilon_k = \frac{eD_T}{\mu_e} \quad (7)$$

The values of electron energy distribution function  $f(\varepsilon)$  are calculated from Boltzmann equation using all the electron collision cross sections as follows,

$$Q_m^T(\varepsilon) = Q_m(\varepsilon) + \sum_j Q_e(\varepsilon) + Q_i(\varepsilon) + Q_a(\varepsilon) \quad (8)$$

Here,  $Q_m^T(\varepsilon)$  represented the total effective momentum transfer cross section, The cross sections  $Q_m(\varepsilon)$ ,  $Q_e(\varepsilon)$ ,  $Q_i(\varepsilon)$  and  $Q_a(\varepsilon)$  are indicate the momentum transfer, excitation (vibration and electronic), ionization, and attachment, respectively.

The reduced ionization coefficient is (Tuan, 2014, 2016),

$$\frac{\alpha}{N} = \frac{1}{v_d} \sqrt{\frac{2e}{m}} \int_i^\infty Q_i(\varepsilon) f(\varepsilon) \varepsilon d\varepsilon \quad (9)$$

The reduced attachment coefficient is,

$$\frac{\eta}{N} = \frac{1}{v_d} \sqrt{\frac{2e}{m}} \int_a^\infty Q_a(\varepsilon) f(\varepsilon) \varepsilon d\varepsilon \quad (10)$$

Where,  $Q_i(\varepsilon)$ ,  $Q_a(\varepsilon)$  are ionization and attachment cross section, here,  $i$  and  $a$  is the ionization and attachment threshold energy. The reduced critical electric field strength  $(E/N)_{cr}$  is calculated when the formation and loss electrons reach a balance, this mean that the effective ionization equal to zero ( Zhao, et al., 2013).

$$\bar{\alpha} = \frac{\alpha}{N} - \frac{\eta}{N} = \frac{\alpha - \eta}{N} = 0 \quad (11)$$

### 3. COLLISION CROSS SECTION

To apply the Boltzmann equation analysis, the electron energy distribution function (EEDF) and values of transport coefficients in nitrous oxide (N<sub>2</sub>O) gas are calculated from the

knowledge of electron collision cross section (elastic and inelastic) which reported by. This sets includes 9 collisional processes: one momentum-transfer cross section ( $Q_m$ ), three vibration excitation ( $Q_{v1}$ ,  $Q_{v2}$  and  $Q_{v3}$ ) with threshold energy 0.073, 0.159, and 0.276 eV respectively and three electronic excitation ( $Q_{ex}$ ) cross section with threshold energy of 4.05, 8.5 and 9.6 eV, one dissociative attachment cross section ( $Q_a$ ), and one ionization cross section ( $Q_i$ ) with threshold energy 12.89 eV.

For  $SF_6$  the cross sections are, momentum transfer cross section  $Q_m$  was taken, while the sets of vibration excitation cross sections  $Q_v$  with threshold energy 0.095 eV was taken from. The ionization cross sections having onset energy 15.8 eV is taken from, electron excitation cross section  $Q_{ex}$  estimated from. The electron attachment cross sections  $Q_a$  are used in the calculation, for low energy  $E \leq 0.140$  eV is estimated the measurements of, the value for  $E > 0.14$  eV is taken from measurements.

#### 4. RESULTS AND DISCUSSION

Many previous papers concentrated on the investigations of EEDF, electron swarm parameters of pure  $N_2O$ , and  $SF_6$ , at temperature 300 K and pressure 1 atm. This paper also calculated the density-reduced effective ionization coefficient  $(\alpha - \eta)/N$  and the density-normalized critical field strength  $(E/N)_{cr}$  of  $N_2O-SF_6$  mixtures using the cross section data reported in section 3.

The influence of the electron distribution function (EEDF) as a function of the mean electron energy for different ratio of reduced electric field ( $E/N$ ) in  $N_2O$  gas, is shown in figure 1. Clearly, the electron energy distribution function is strongly affected by changing the parameter  $E/N$ , subsequently the electron transport coefficients depends on the ratio  $E/N$ . Near the thresholds of the inelastic processes, the formation of the EEDF is strongly influenced by the electric field, which heats the electron component and thus increases the energy of the cold electrons. For electron energies in the range less than 2 eV, the EEDF decreases as the  $E/N$  increases, however for electron energies greater than 2 eV the EEDF increases with increasing  $E/N$ , because at high value of  $E/N$  the kinetic energy of electrons increases. It is obviously appears from the figure that applying a high electric field (or  $E/N$ ) leads to the development of EEDF to a higher energy tail.

Figure 2 is shown the influence of superelastic collisions on EEDF at reduced electric field strength 0.5, 5 and 30 Td with and without superelastic collision. It can be seen that superelastic collisions are effect the EEDF at low reduced electric field strength  $E/N$ , but are not important at high  $E/N$  values. For high values of reduced electric field strength  $E/N$ , the electrons gain their energy from the applied d.c electric field. The study of the influence of second kind (superelastic) collisions was explained in literatures of (Pietanza, et al., 2016; Pietanza, et al., 2016a).

Figure 3 shows the EEDF in pure  $N_2O$ ,  $SF_6$  and 50%  $N_2O$ -50%  $SF_6$  gas mixtures under gas temperature of 300 K and  $E/N$  of 100 Td. It can be seen from the figure that the number of electrons in the low-energy region (electron energy less than 2 eV) increased but the number of high-energy electrons decreased. This is mainly due to the fact that  $N_2O$  is kind of strong electronegative gas, at low energy region larger amount of electrons were adsorbed and accumulated in pure  $N_2O$  gas.

Figure 4 shows the mean electron energy from 0.025 eV at  $E/N=0.1$  Td to 13.5 eV at  $E/N=1000$  Td, at high  $E/N$  the mean electron energy sensitive to inelastic collisions. The variation of mean electron energy is progressive exponential, the electron gain all energies from the applied electric field. Comparison has been made with the theoretical values of and experimental values of (Mechlinska-Drewko, et al., 2003; Lisovski, et al., 2006) a good

agreement has been observed. The results of electron drift velocity in  $N_2O$  as a function of  $E/N$  is shown in figure 5, which increases with increasing  $E/N$  values, the results are compared with the theoretical values of and experimental values of (Dupljanin, et al., 2010; Basurto, et al., 2013). The agreement is good over the entire reduced electric field strength  $E/N$  range, on the other hand the theoretical results of, and experimental results of (Basurto, et al., 2013) are little higher than present calculation. Figure 6 shows the calculated results for the electron drift velocity in  $SF_6$  as a function of  $E/N$ , in comparison with the experimental data of (Aschwanden, 1985; Lisovski, et al., 2010) as well as the theoretical results of (Wei, et al., 2014; Miric, et al., 2016; Seasa, et al., 2019), good agreements has been observed. On the other hand the experimental results of (Lisovski, et al., 2010), throughout the whole range of  $90Td \leq E/N \leq 200Td$  is higher than the present calculation. Figure 7 gives the electron drift velocities, calculated in this work for 100%  $N_2O$ , 100%  $SF_6$  and 50%  $N_2O$ -50%  $SF_6$  mixture as a function of the  $E/N$ . It can be clearly observed at fixed value of  $E/N$  that the electron drift velocity displays a quick trend to increase as the  $N_2O$  content increases, especially in the low  $E/N$  range.

The characteristic energy in pure  $N_2O$  as a function of  $E/N$  is shown in figure 8. The present calculation were found in good agreement with experimental results of (Mechlinska-Drewko, 2003; Dupljanin, et al., 2010).

Figure 9, shows the calculated results for the normalized reduced ionization coefficient  $\alpha/N$  as a function of  $E/N$ , in comparison with the experimental data of (Lisovski, et al., 2006) as well as the theoretical results of (Dupljanin, et al., 2010). Throughout the whole range of  $70Td \leq E/N \leq 1000Td$  good agreements has been observed. For  $E/N < 350Td$ , the experimental results of lower compare with the present results. In addition, the experimental results of (Lisovski, et al., 2006) over the range of  $E/N < 250Td$  higher than present results. Figure 10 shows the comparison of the calculated normalized density attachment coefficient  $\eta/N$  in pure  $N_2O$  with the previous literatures (Lisovski, et al., 2006; Dupljanin, et al., 2010). A deviation between the calculated and experimental data of (Yoshida, et al., 1999; Lisovski, et al., 2006) can be observed due to the cross-sections for the attachment, However, the calculated results show good agreement with theoretical results of (Dupljanin, et al., 2010) at room temperature, and experimental results.

In the present study, the results from the two-term approximation Boltzmann equation analysis and the data from the experimental measurements (Basurto, et al., 2007; Basurto, et al., 2013) and theoretical calculation (Dupljanin, et al., 2010) of the density-reduced effective ionization coefficient in pure  $N_2O$  are compared in Fig. 11. It can be observed that the density normalized coefficient  $(\alpha - \eta)/N$  increases with increasingly reducing the field strength. Although the experimental results of is lower than present calculation over the range of  $E/N < 90Td$ . The density-normalized critical field strength  $(E/N)_{cr}$  values at which ionization coefficient balance with attachment coefficient ( $\alpha/N = \eta/N$ ) is obtained in the present calculation is equal to approximately 199Td.

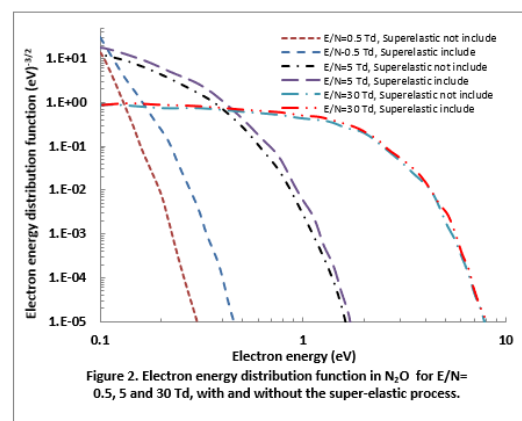
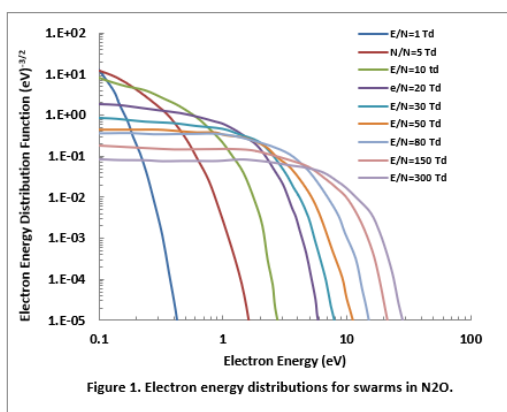
In figure 12 we show variation of the gas density reduced ionization  $\alpha/N$  and attachment coefficient  $\eta/N$  with  $E/N$  for electrons in  $SF_6$ , by using the two-term Boltzmann equation analysis, which theoretical and experimental values for comparison are mention in the figure. It is seen that the ionization coefficient increase and attachment coefficient decreases as reduced electric field strength  $E/N$  increases. The attachment coefficient is defined the probability that an electron will attach to a gas molecule in traveling a unit distance in the electric field direction is a function of  $E/N$ . The agreement between the present calculation and previous experimental values of, and theoretical results of (Abderrahmane and Lahouaria, 2005; Miric, et al., 2016) are excellent for ionization coefficient, on the other hand the

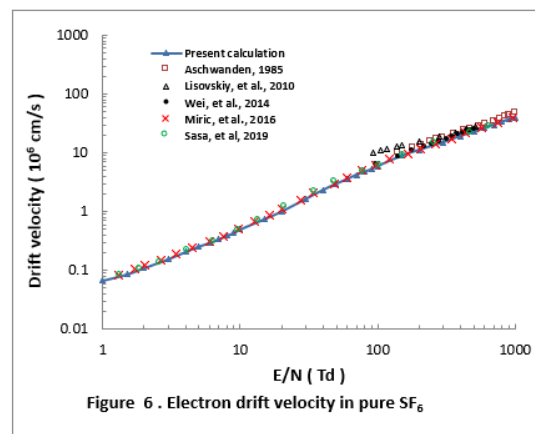
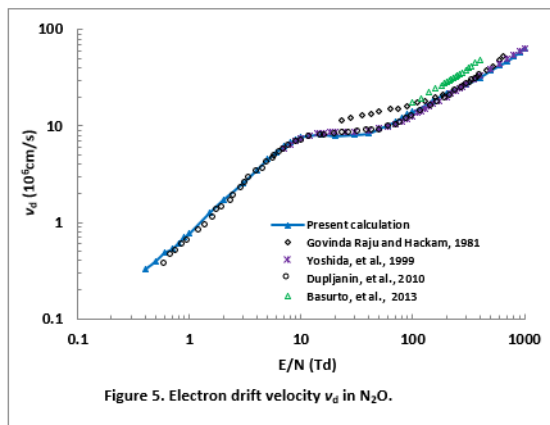
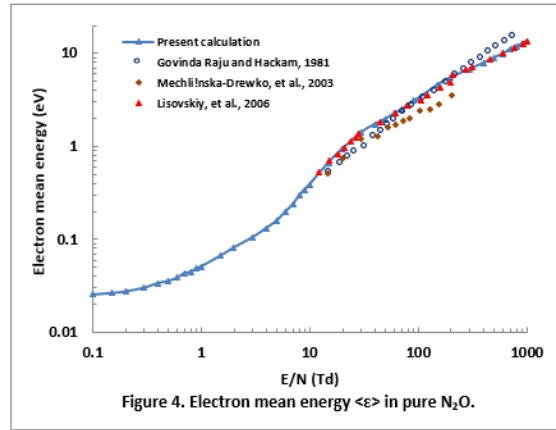
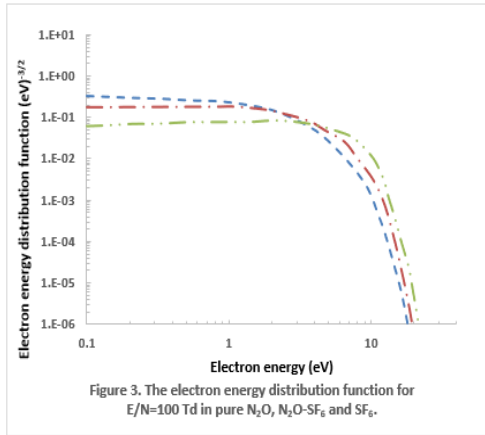
theoretical results of Itoh higher than that of present results for the range of  $E/N < 215$  Td by %7.

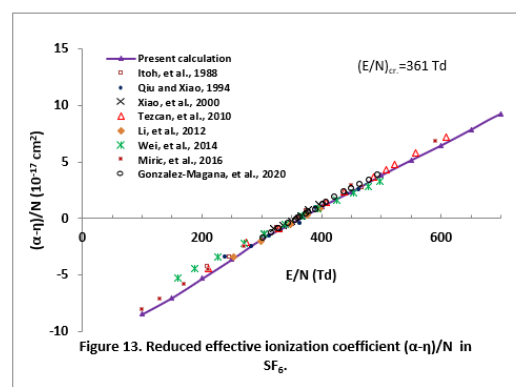
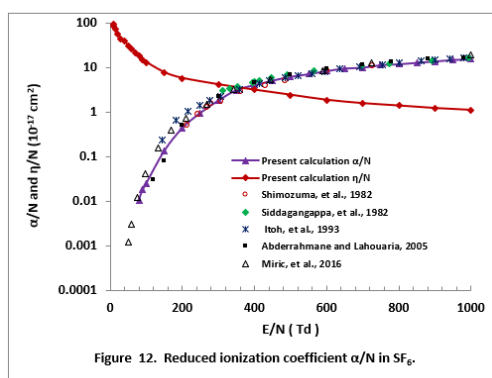
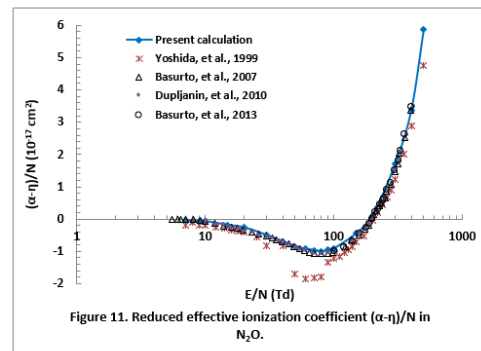
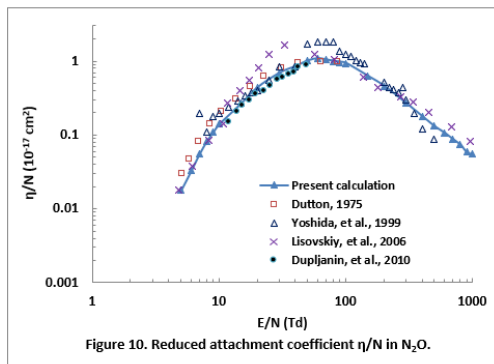
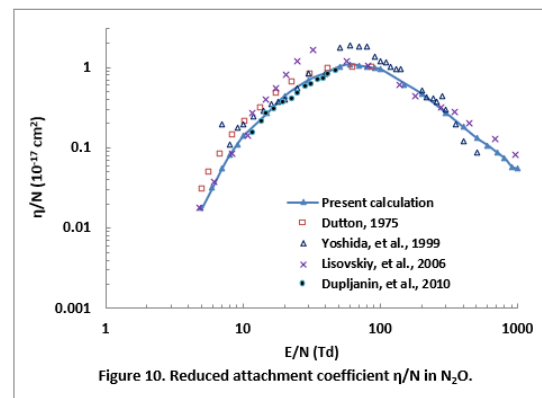
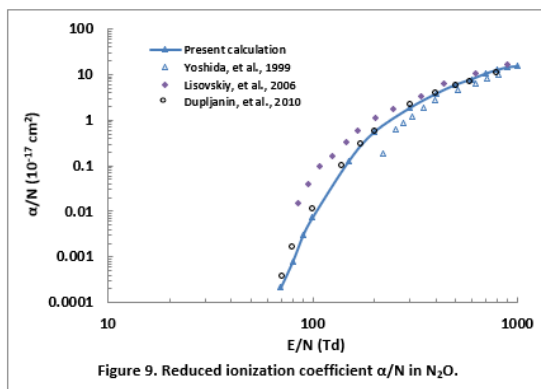
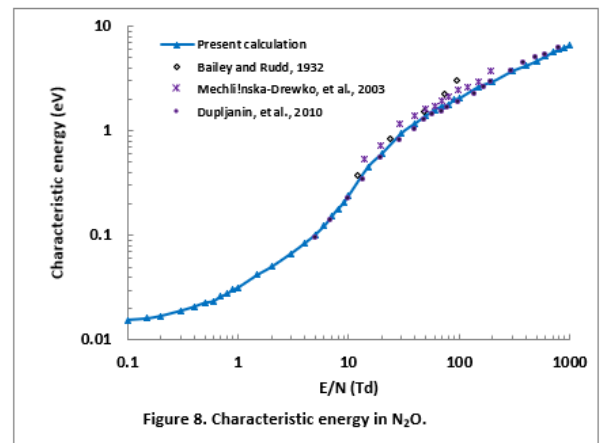
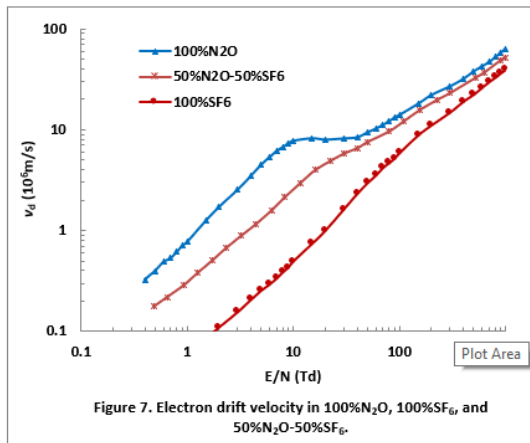
A zero value of  $(\alpha - \eta)/N$  means a critical breakdown voltage of a gas, and the critical field strength  $(E/N)_{cr}$  of the pure and binary mixture was calculated as the value of reduced electric field strength at  $(\alpha - \eta)/N = 0$ , and the critical electric field strength  $E_{cr}$  was then calculated by multiplying  $(E/N)_{cr}$  by  $N$ , where  $N$  is the gas density of the gas. Figure 13 shows the density normalized effective ionization coefficient  $(\alpha - \eta)/N$  in  $SF_6$  calculated over a range of  $E/N$  values from 100 Td to 700 Td by using two-term solution of Boltzmann equation. The critical field strength  $(E/N)_{cr}$  values at which  $(\alpha - \eta)/N = 0$  is obtained in the present calculation is equal to 361 Td, in agreement with the experimental results of: 361 Td, 361 Td (Xiao et al., 2000), 361 Td (Li, et al., 2012), 360 Td Gonzalez-Magaña, et al., 2020), and theoretical results of: 362 Td, 360 (Tezcan, et al., 2010), 361 (Wei, et al., 2014), 361 (Mirić, et al., 2016). The effective ionization coefficient is almost zero, since ionizing collisions are balanced by attaching collisions, and for  $E/N$  values smaller than the  $(E/N)_{cr}$ , attachment processes become dominant, yielding negative values for the effective ionization coefficient as  $E/N$  is decreased and; on the other hand, for  $E/N$  values above the  $(E/N)_{cr}$ , the effective ionization coefficient increases with increasing  $E/N$  values where the ionization collisions become dominant and the effect of the attachment processes is not significant.

The linear relationship between the density-reduced effective ionization coefficient  $(\alpha - \eta)/N$  and reduced electric field strength  $E/N$  is common for electronegative gases, such as  $SF_6$  and  $SF_6$  mixtures (Christophorou and Brunt, 2002). The law is verified valid for  $N_2O-SF_6$  mixtures it can be represented by,  $[(\alpha - \eta)/N](10^{-17} \text{ cm}^2) = \beta[(E/N) - (E/N)_{cr}]$ , where  $(E/N)_{cr}$  is the critical field strength value, and  $\beta$  is the slope of the curve  $(\alpha - \eta)/N = f(E/N)$ . The values of  $\beta$  and  $(E/N)_{cr}$  depend on the content of  $SF_6$  in the mixtures, which are shown in figure 14. The density-reduced effective ionization coefficient  $(\alpha - \eta)/N$  increase with increasing reduced field strength  $E/N$ , and the critical field strength  $(E/N)_{cr}$  of  $N_2O-SF_6$  mixture increase with increasing  $SF_6$  molar fraction in the gas mixture.

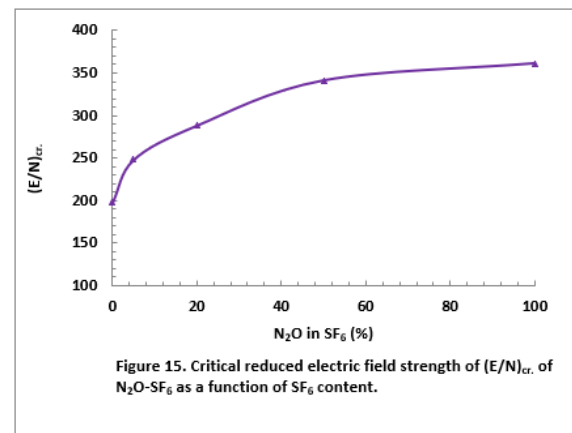
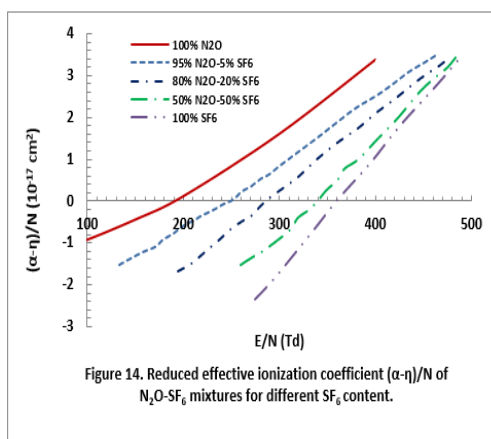
The values of density-normalized critical electric field strength  $(E/N)_{cr}$  are presented in Figure 15 as a function of  $N_2O$  content in the binary mixtures of  $SF_6$  as a parameter with concentrations of 5%, 20% and 50%. The density-normalized critical electric field strength of the binary mixture  $N_2O-SF_6$  increases with increasing concentration of  $SF_6$ .











## 5. CONCLUSION

The electron swarm parameters, namely (drift velocity, electron mean energy, characteristic energy, ionization and attachment coefficients), density-reduced effective ionization coefficient  $(\alpha - \eta)/N$  and density-normalized critical field strength in pure  $N_2O$ ,  $SF_6$  and binary  $N_2O-SF_6$  mixtures have been calculated using two term spherical harmonic approximation of the Boltzmann equation analysis at temperature 300 K and pressure 1 atm. The overall  $E/N$  is from 0.1 to 1000 Td, while the  $N_2O$  content in the binary  $N_2O-SF_6$  mixtures can be varied from 100% to 0%. From the zeroth value of density-reduced effective ionization coefficient  $(\alpha - \eta)/N$ , the density-normalized critical field strength  $(E/N)_{cr}$  for pure gases  $N_2O$  (199 Td),  $SF_6$  (361 Td), and for each  $N_2O$  concentration is derived. The critical field strength increase with increasing  $SF_6$  concentration in the gas mixtures.

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